

Cluster algebra and generalized q -Painlevé VI systems of type A

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Higher order generalizations of the q -Painlevé equations (we call it q -Painlevé systems) has been studied from some points of view.

- Birational representations of extended affine Weyl groups
 - Type $(A_{m-1} + A_{n-1})^{(1)}$ [Kajiwara-Noumi-Yamada 02]
 - Type $D_{2n+3}^{(1)}$ (q -Sasano system) [Masuda 15]
- Cluster mutations
 - Somos-type Y -system [Hone-Inoue 14]
- Compatibility conditions of Lax pairs
 - q -Garnier system [Sakai 05], [Nagao-Yamada 18]
 - Similarity reduction of lattice q -UC hierarchy [Tsuda 10]
 - Similarity reduction of q -Drinfeld-Sokolov hierarchy [S 15], [S 17]

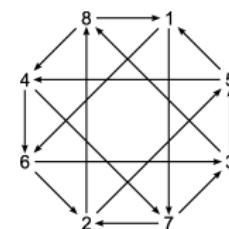
In this talk we present the following result.

- ① A birational representation of $\widetilde{W}((A_{2n+1} + A_1 + A_1)^{(1)})$ is formulated with the aid of cluster mutations.
- ② This extended affine Weyl group provides the latter three q -Painlevé systems (arising from Lax pairs) as translations.

We recall the derivation of Jimbo-Sakai's q - P_{VI} from cluster mutations.

Consider a skew-symmetric matrix corresponding to the q - P_{VI} quiver (cf. [Okubo 15])

$$\Lambda = (\lambda_{i,j})_{i,j=1}^8 = \begin{pmatrix} 0 & 0 & 0 & 0 & -1 & 1 & 1 & -1 \\ 0 & 0 & 0 & 0 & 1 & -1 & -1 & 1 \\ 0 & 0 & 0 & 0 & 1 & -1 & -1 & 1 \\ 0 & 0 & 0 & 0 & -1 & 1 & 1 & -1 \\ 1 & -1 & -1 & 1 & 0 & 0 & 0 & 0 \\ -1 & 1 & 1 & -1 & 0 & 0 & 0 & 0 \\ -1 & 1 & 1 & -1 & 0 & 0 & 0 & 0 \\ 1 & -1 & -1 & 1 & 0 & 0 & 0 & 0 \end{pmatrix}.$$



Example

A mutation $\mu_1 : (\Lambda, \mathbf{y}) \mapsto (\Lambda', \mathbf{y}')$ is given explicitly by

$$\begin{aligned} \lambda'_{5,6} &= \lambda'_{5,7} = \lambda'_{8,6} = \lambda'_{8,7} = 1, & \lambda'_{6,5} &= \lambda'_{7,5} = \lambda'_{6,8} = \lambda'_{7,8} = -1, \\ \lambda'_{1,i} &= -\lambda_{1,i}, & \lambda'_{i,1} &= -\lambda_{i,1} \quad (i \neq 1), & \lambda'_{i,j} &= \lambda_{i,j} \quad (\text{otherwise}), \\ y'_5 &= y_5(1 + y_1), & y'_6 &= \frac{y_6}{1 + \frac{1}{y_1}}, & y'_7 &= \frac{y_7}{1 + \frac{1}{y_1}}, & y'_8 &= y_8(1 + y_1), \\ y'_1 &= \frac{1}{y_1}, & y'_i &= y_i \quad (\text{otherwise}). \end{aligned}$$

We can formulate a birational representation of $\widetilde{W}(A_3^{(1)})$ with the aid of cluster mutations.

Fact (cf. [Bershtein-Gavrylenko-Marshakov 17])

Let

$$\begin{aligned} r_0 &= (1, 2) \mu_1 \mu_2, & r_1 &= (5, 6) \mu_5 \mu_6, & r_2 &= (3, 4) \mu_3 \mu_4, & r_2 &= (7, 8) \mu_7 \mu_8, \\ \pi &= (1, 5, 3, 7)(2, 6, 4, 8), \end{aligned}$$

where μ_k stands for a mutation at a vertex k and (i, j) a permutation of vertices i and j . Then the q - P_{VI} quiver is invariant under their actions. Moreover they satisfy

$$r_i^2 = 1, \quad (r_i r_{i+1})^3 = 1, \quad (r_i r_j)^2 = 1 \quad (j \neq i, i \pm 1), \quad \pi^4 = 1, \quad r_i \pi = \pi r_{i+1},$$

where $r_{i+4} = r_i$.

Fact (cf. [Okubo 15])

A translation $T = (\pi r_0 r_2)^2$ of the extended affine Weyl group provides q - P_{VI} .

In this talk we consider an extension of this previous work.

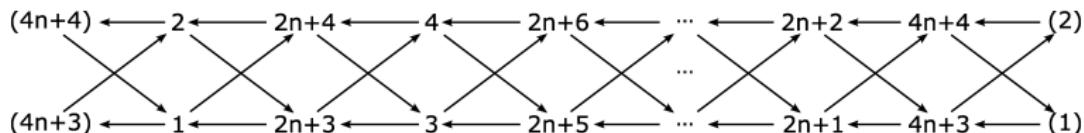
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Consider the following quiver on a torus (cf. [Inoue-Ishibashi-Ohya 18]).



Let y_1, \dots, y_{4n+4} be coefficients. We define parameters by

$$\alpha_{2i} = y_{2i+1} y_{2i+2}, \quad \alpha_{2i+1} = y_{2i+2n+3} y_{2i+2n+4} \quad (i = 0, \dots, n), \quad q = \prod_{i=0}^{2n+1} \alpha_i,$$

$$\beta = \prod_{i=0}^n \frac{1}{y_{2i+2} y_{2i+2n+3}}, \quad \beta' = \prod_{i=0}^n y_{2i+1} y_{2i+2n+3},$$

and dependent variables by

$$\varphi_{2i} = y_{2i+1}, \quad \varphi_{2i+1} = y_{2i+2n+3} \quad (i = 0, \dots, n).$$

We use a notation for a periodicity

$$\alpha_{i+2n+2} = \alpha_i, \quad \varphi_{i+2n+2} = \varphi_i.$$

Definition (cf. [Bershtein-Gavrylenko-Marshakov 17])

We define reflections r_0, \dots, r_{2n+1} by

$$\begin{aligned}r_{2i} &= (2i+1, 2i+2) \mu_{2i+2} \mu_{2i+1}, \\r_{2i+1} &= (2i+2n+3, 2i+2n+4) \mu_{2i+2n+4} \mu_{2i+2n+3}.\end{aligned}$$

They act on $(\alpha_i, \beta, \beta'; \varphi_i)$ as

$$\begin{aligned}r_i(\alpha_{i-1}) &= \alpha_i \alpha_{i-1}, & r_i(\alpha_i) &= \frac{1}{\alpha_i}, & r_i(\alpha_{i+1}) &= \alpha_i \alpha_{i+1}, & r_i(\alpha_j) &= \alpha_j, \\r_i(\beta) &= \beta, & r_i(\beta') &= \beta', \\r_i(\varphi_{i-1}) &= \frac{\alpha_i + \varphi_i}{1 + \varphi_i} \varphi_{i-1}, & r_i(\varphi_i) &= \frac{\varphi_i}{\alpha_i}, & r_i(\varphi_{i+1}) &= \frac{1 + \varphi_i}{1 + \frac{\varphi_i}{\alpha_i}} \varphi_{i+1}, & r_i(\varphi_j) &= \varphi_j,\end{aligned}$$

for $j \neq i, i \pm 1$ (cf. [Kajiwara-Noumi-Yamada 02]).

Definition (cf. [Inoue-Lam-Polyavskyy 16], [Masuda-Okubo-Tsuda 18])

We define reflections s_0, s_1, s'_0, s'_1 by

$$s_0 = \mu_1 \mu_{2n+4} \mu_3 \mu_{2n+6} \cdots \mu_{2n-1} \mu_{4n+2} \mu_{2n+1} (2n+1, 4n+4)$$

$$\mu_{2n+1} \mu_{4n+2} \mu_{2n-1} \cdots \mu_{2n+6} \mu_3 \mu_{2n+4} \mu_1,$$

$$s_1 = \mu_2 \mu_{2n+3} \mu_4 \mu_{2n+5} \cdots \mu_{2n} \mu_{4n+1} \mu_{2n+2} (2n+2, 4n+3)$$

$$\mu_{2n+2} \mu_{4n+1} \mu_{2n} \cdots \mu_{2n+5} \mu_4 \mu_{2n+3} \mu_2,$$

$$s'_0 = \mu_1 \mu_{2n+3} \mu_3 \mu_{2n+5} \cdots \mu_{2n-1} \mu_{4n+1} \mu_{2n+1} (2n+1, 4n+3)$$

$$\mu_{2n+1} \mu_{4n+1} \mu_{2n-1} \cdots \mu_{2n+5} \mu_3 \mu_{2n+3} \mu_1,$$

$$s'_1 = \mu_2 \mu_{2n+4} \mu_4 \mu_{2n+6} \cdots \mu_{2n} \mu_{4n+2} \mu_{2n+2} (2n+2, 4n+4)$$

$$\mu_{2n+2} \mu_{4n+2} \mu_{2n} \cdots \mu_{2n+6} \mu_4 \mu_{2n+4} \mu_2.$$

They act on $(\alpha_i, \beta, \beta')$ as

$$s_k(\alpha_i) = \alpha_i, \quad s_k(\beta') = \beta' \quad (k = 0, 1), \quad s_0(\beta) = \frac{1}{q^2 \beta}, \quad s_1(\beta) = \frac{1}{\beta},$$

$$s'_l(\alpha_i) = \alpha_i, \quad s'_l(\beta) = \beta \quad (l = 0, 1), \quad s'_0(\beta') = \frac{1}{\beta'}, \quad s'_1(\beta') = \frac{q^2}{\beta'}.$$

They also act on (φ_i) as

$$\begin{aligned}
 s_0(\varphi_{2i}) &= \frac{\varphi_{2i+1}}{\alpha_{2i+1}} \frac{\sum_{j=0}^n (\prod_{k=0}^{j-1} \varphi_{2i+2k} \frac{\alpha_{2i+2k+1}}{\varphi_{2i+2k+1}}) (1 + \varphi_{2i+2j})}{\sum_{j=0}^n (\prod_{k=0}^{j-1} \varphi_{2i+2k+2} \frac{\alpha_{2i+2k+3}}{\varphi_{2i+2k+3}}) (1 + \varphi_{2i+2j+2})}, \\
 s_0(\varphi_{2i+1}) &= \alpha_{2i+1} \varphi_{2i+2} \frac{\sum_{j=0}^n (\prod_{k=0}^{j-1} \frac{\alpha_{2i+2k+3}}{\varphi_{2i+2k+3}} \varphi_{2i+2k+4}) (1 + \frac{\alpha_{2i+2j+3}}{\varphi_{2i+2j+3}})}{\sum_{j=0}^n (\prod_{k=0}^{j-1} \frac{\alpha_{2i+2k+1}}{\varphi_{2i+2k+1}} \varphi_{2i+2k+2}) (1 + \frac{\alpha_{2i+2j+1}}{\varphi_{2i+2j+1}})}, \\
 s_1(\varphi_{2i}) &= \alpha_{2i} \varphi_{2i+1} \frac{\sum_{j=0}^n (\prod_{k=0}^{j-1} \frac{\alpha_{2i+2k+2}}{\varphi_{2i+2k+2}} \varphi_{2i+2k+3}) (1 + \frac{\alpha_{2i+2j+2}}{\varphi_{2i+2j+2}})}{\sum_{j=0}^n (\prod_{k=0}^{j-1} \frac{\alpha_{2i+2k}}{\varphi_{2i+2k}} \varphi_{2i+2k+1}) (1 + \frac{\alpha_{2i+2j}}{\varphi_{2i+2j}})}, \\
 s_1(\varphi_{2i+1}) &= \frac{\varphi_{2i+2}}{\alpha_{2i+2}} \frac{\sum_{j=0}^n (\prod_{k=0}^{j-1} \varphi_{2i+2k+1} \frac{\alpha_{2i+2k+2}}{\varphi_{2i+2k+2}}) (1 + \varphi_{2i+2j+1})}{\sum_{j=0}^n (\prod_{k=0}^{j-1} \varphi_{2i+2k+3} \frac{\alpha_{2i+2k+4}}{\varphi_{2i+2k+4}}) (1 + \varphi_{2i+2j+3})}, \\
 s'_0(\varphi_i) &= \frac{1}{\varphi_{i+1}} \frac{\sum_{j=0}^{2n+1} \prod_{k=0}^{j-1} \frac{1}{\varphi_{i+k+2}}}{\sum_{j=0}^{2n+1} \prod_{k=0}^{j-1} \frac{1}{\varphi_{i+k}}}, \\
 s'_1(\varphi_i) &= \frac{\alpha_i}{\frac{\varphi_{i+1}}{\alpha_{i+1}}} \frac{\sum_{j=0}^{2n+1} \prod_{k=0}^{j-1} \frac{\varphi_{i+k}}{\alpha_{i+k}}}{\sum_{j=0}^{2n+1} \prod_{k=0}^{j-1} \frac{\varphi_{i+k+2}}{\alpha_{i+k+2}}}.
 \end{aligned}$$

Definition

We define rotations π, π' by

$$\begin{aligned}\pi = & (1, 2n+3, 3, 2n+5, \dots, 2n+1, 4n+3, 1) \\ & (2, 2n+4, 4, 2n+6, \dots, 2n+2, 4n+4, 2), \\ \pi' = & (1, 2n+4, 3, 2n+6, \dots, 2n+1, 4n+4, 1) \\ & (2, 2n+3, 4, 2n+5, \dots, 2n+2, 4n+3, 2).\end{aligned}$$

They act on $(\alpha_i, \beta, \beta'; \varphi_i)$ as

$$\begin{aligned}\pi(\alpha_i) &= \alpha_{i+1}, \quad \pi(\beta) = \frac{1}{q\beta}, \quad \pi(\beta') = \beta', \quad \pi(\varphi_i) = \varphi_{i+1}, \\ \pi'(\alpha_i) &= \alpha_{i+1}, \quad \pi'(\beta) = \beta, \quad \pi'(\beta') = \frac{q}{\beta'}, \quad \pi'(\varphi_i) = \frac{\alpha_{i+1}}{\varphi_{i+1}}.\end{aligned}$$

Theorem ([Okubo-S 18], [Masuda-Okubo-Tsuda 18])

The birational transformations $r_0, \dots, r_{2n+1}, s_0, s_1, s'_0, s'_1, \pi, \pi'$ satisfy the fundamental relations of an extended affine Weyl group of type $(A_{2n+1} + A_1 + A_1)^{(1)}$

$$r_i^2 = 1, \quad (r_i r_{i+1})^3 = 1, \quad (r_i r_j)^2 = 1 \quad (j \neq i, i \pm 1),$$

$$s_k^2 = 1, \quad (s_l')^2 = 1, \quad (s_k s_l')^2 = 1,$$

$$\pi^{2n+2} = 1, \quad (\pi')^{2n+2} = 1, \quad \pi\pi' = \pi'\pi, \quad \pi^2 = (\pi')^2,$$

$$(r_i s_k)^2 = 1, \quad (r_i s_l')^2 = 1,$$

$$r_i \pi = \pi r_{i+1}, \quad r_i \pi' = \pi' r_{i+1},$$

$$s_k \pi = \pi s_{k+1}, \quad s_k \pi' = \pi' s_k, \quad s_l' \pi = \pi s_l', \quad s_l' \pi' = \pi' s_{l+1}',$$

where

$$r_{i+2n+2} = r_i, \quad s_{k+2} = s_k, \quad s_{l+2}' = s_l'.$$

Remark

The fundamental relation $(s_k s_l')^2 = 1$ was shown by R. Inoue, T. Lam and P. Pylyavskyy in 2016.

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Definition

We define translations T_1, T_2, T_3 and a quasi-translation T_4 by

$$\begin{aligned} T_1 &= s'_1 s_1 \pi' \pi^{-1}, & T_2 &= (r_2 r_3 \dots r_{2n+1} \pi')^2, & T_3 &= r_1 r_2 \dots r_{2n+1} s'_1 \pi', \\ T_4 &= (r_0 r_2 \dots r_{2n} \pi')^2. \end{aligned}$$

They act on $(\alpha_i, \beta, \beta')$ as

$$\begin{aligned} T_1(\beta) &= q \beta, & T_1(\beta') &= q \beta', \\ T_2(\alpha_0) &= \frac{\alpha_0}{q}, & T_2(\alpha_n) &= q \alpha_n, & T_2(\alpha_{n+1}) &= \frac{\alpha_{n+1}}{q}, & T_2(\alpha_{2n+1}) &= q \alpha_{2n+1}, \\ T_3(\alpha_0) &= q \alpha_0, & T_3(\alpha_1) &= \frac{\alpha_1}{q}, & T_3(\beta) &= q \beta, \\ T_4(\alpha_{2i}) &= \frac{1}{\alpha_{2i+1} \alpha_{2i+2} \alpha_{2i+3}}, & T_4(\alpha_{2i+1}) &= \alpha_{2i+1} \alpha_{2i+2} \alpha_{2i+3} \alpha_{2i+4} \alpha_{2i+5}. \end{aligned}$$

Note that iterative action of T_4 becomes a translation as

$$T_4^m(\alpha_{2i}) = \frac{\alpha_{2i}}{q}, \quad T_4^m(\alpha_{2i+1}) = q \alpha_{2i+1} \quad (n = 2m - 1),$$

$$T_4^{2m+1}(\alpha_{2i}) = \frac{\alpha_{2i}}{q^2}, \quad T_4^{2m+1}(\alpha_{2i+1}) = q^2 \alpha_{2i+1} \quad (n = 2m).$$

Theorem ([Okubo-S 18])

These translations gives the known three higher order q -Painlevé systems.

- T_1 : q -DS system $q\text{-}P_{(n+1,n+1)}$
- T_2 : *Sakai's q -Garnier system.*
- T_4 : *Tsuda's q -LUC system*

Conjecture

- T_3 : *Nagao-Yamada's variation of the q -Garnier system*

T_1 : q -DS system $q\text{-}P_{(n+1,n+1)}$

The following lemma is obtained in the course of proof of the fundamental relations.

Lemma

We have

$$s'_1 s_1(\varphi_{2i-1}) = s_1 s'_1(\varphi_{2i-1}) = \frac{1 + \frac{\alpha_{2i}}{\varphi_{2i}}}{1 + \frac{\alpha_{2i-2}}{\varphi_{2i-2}}} \frac{\alpha_{2i+1}}{\varphi_{2i+1}} \frac{S_{2i-2} S'_{2i} + \frac{\alpha_{2i-2}}{\varphi_{2i-2}} \alpha_{2i-1} S'_{2i-2} S_{2i}}{S_{2i} S'_{2i+2} + \frac{\alpha_{2i}}{\varphi_{2i}} \alpha_{2i+1} S'_{2i} S_{2i+2}},$$

$$s'_1 s_1(\varphi_{2i}) = s_1 s'_1(\varphi_{2i}) = \frac{\alpha_{2i} \alpha_{2i+1}}{\frac{\varphi_{2i}}{\alpha_{2i}}} \frac{S_{2i+2} S'_{2i}}{S_{2i} S'_{2i+2}},$$

where

$$S_{2i} = \sum_{j=0}^n \left(\prod_{k=0}^{j-1} \frac{\alpha_{2i+2k}}{\varphi_{2i+2k}} \varphi_{2i+2k+1} \right) \left(1 + \frac{\alpha_{2i+2j}}{\varphi_{2i+2j}} \right),$$

$$S'_{2i} = \sum_{j=0}^n \left(\prod_{k=0}^{j-1} \frac{\varphi_{2i+2k}}{\alpha_{2i+2k}} \frac{\varphi_{2i+2k+1}}{\alpha_{2i+2k+1}} \right) \left(1 + \frac{\varphi_{2i+2j}}{\alpha_{2i+2j}} \right).$$

T_1 : q -DS system $q\text{-}P_{(n+1,n+1)}$

Let a_i, b_i ($i = 1, \dots, n+1$) be parameters and t an independent variable defined by

$$\frac{a_i}{b_i} = \alpha_{2i-1}, \quad \frac{b_i}{a_{i+1}} = \alpha_{2i}, \quad q^{n-1} t \prod_{i=1}^{n+1} a_i b_i = \beta, \quad t = \beta'.$$

Also let f_i, g_i ($i = 0, \dots, n$) be dependent variables defined by

$$f_i = \frac{1 + \varphi_{2i}}{1 + \varphi_0} \prod_{j=i}^n \varphi_{2j+1} \varphi_{2j+2}, \quad \frac{g_i}{b_i} = -\frac{1}{\varphi_{2i}}.$$

Note that

$$f_0 = t, \quad g_0 = \frac{1}{q^{\frac{n-2}{2}} t g_1 \dots g_n}, \quad b_0 = q b_{n+1}.$$

T_1 : q -DS system $q\text{-}P_{(n+1,n+1)}$

Theorem ([Okubo-S 18])

If we set $\bar{f}_i = T_1(f_i)$ and $\bar{g}_i = T_1(g_i)$, then they satisfy

$$f_i \bar{f}_i = q t \frac{F_i F_{i+1} \left(\frac{b_i}{\bar{g}_i} - 1 \right) (\bar{g}_i - a_{i+1})}{F_{n+1} F_1 \left(\frac{b_0}{\bar{g}_0} - 1 \right) (\bar{g}_0 - a_1)}, \quad g_i \bar{g}_i = \frac{F_{i+1} G_i}{F_i G_{i+1}} \quad (i = 1, \dots, n),$$

where

$$F_i = \sum_{j=1}^{i-1} f_j + t \sum_{j=i}^n f_j + t,$$

$$\begin{aligned} G_i &= \sum_{j=i}^n \prod_{k=i}^{j-1} b_k a_{k+1} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j + q^{\frac{n}{2}} t \prod_{k=i}^n b_k a_{k+1} \\ &\quad + q^n t \sum_{j=1}^{i-1} \frac{b_{n+1} a_1 \prod_{k=1}^n b_k a_{k+1}}{\prod_{k=j}^{i-1} b_k a_{k+1}} \frac{\prod_{l=j+1}^n g_l}{\prod_{l=1}^{j-1} g_l} f_j. \end{aligned}$$

T_2 : q -Garnier system

Fact ([S 2017])

The system q - $P_{(n+1,n+1)}$ is derived from a Lax pair

$$T_{q,z}^{-1}(\psi) = M\psi, \quad T_{q,t}^{-1}(\psi) = B\psi.$$

Moreover, the action of T_2 on (f_i, g_i) is derived from the compatibility condition of

$$(1) \quad T_{q,z}^{-1}(\psi) = M\psi, \quad T_2(\psi) = \Gamma\psi.$$

Here the matrices M , B and Γ are expressed as

$$X = \begin{pmatrix} X_{1,1} & X_{1,2} & & & & \\ & X_{2,2} & X_{2,3} & & & \\ & & X_{3,3} & & & \\ & & & \ddots & & \\ & & & & X_{n,n} & X_{n,n+1} \\ & z X_{n+1,1} & & & & X_{n+1,n+1} \end{pmatrix}, \quad (X = M, B, \Gamma),$$

where each block is a 2×2 matrix.

T_2 : q -Garnier system

Via a q -Laplace transformation $(z, T_{q,z}^{-1}) \rightarrow (T_{q,z}, z)$, system (1) is reduced to the one with 2×2 matrices

$$(2) \quad T_{q,z}(\Psi) = \mathcal{A}\Psi, \quad T_2(\Psi) = \mathcal{B}\Psi,$$

with

$$\begin{aligned} \mathcal{A} &= M_{n+1,1}^{-1}(zI - M_{n+1,n+1}) M_{n,n+1}^{-1}(zI - M_{n,n}) \dots M_{1,2}^{-1}(zI - M_{1,1}), \\ \mathcal{B} &= \Gamma_{1,1} + \Gamma_{1,2} M_{1,2}^{-1}(zI - M_{1,1}). \end{aligned}$$

Lemma

The compatibility condition of system (1) implies the one of (2).

Remark

The transformation from system (1) to (2) is suggested by Y. Yamada.

T_2 : q -Garnier system

System (2) is equivalent to the Lax pair of the q -Garnier system (of inverse direction).

Theorem ([Okubo-S 18])

The matrices \mathcal{A} and \mathcal{B} satisfy the following properties.

$$\textcircled{1} \quad \mathcal{A} = \mathcal{A}_0 + z \mathcal{A}_1 + \dots + z^{n+1} \mathcal{A}_{n+1},$$

$$\mathcal{A}_{n+1} = (-a_1)^{n+1} \begin{pmatrix} t^{-1} & 0 \\ * & 1 \end{pmatrix}, \quad \mathcal{A}_0 \sim \begin{pmatrix} q^{-\frac{n}{2}} t^{-1} & 0 \\ 0 & q^{\frac{n}{2}} a_1 b_1 \dots a_{n+1} b_{n+1} \end{pmatrix}.$$

$$\textcircled{2} \quad \det \mathcal{A} = \frac{a_1^{2n+2}}{t} (z-1) \left(z - \frac{b_1}{a_1}\right) \left(z - \frac{a_2}{a_1}\right) \left(z - \frac{b_2}{a_1}\right) \dots \left(z - \frac{a_{n+1}}{a_1}\right) \left(z - \frac{b_{n+1}}{a_1}\right).$$

$$\textcircled{3} \quad \det \mathcal{B} = \begin{cases} t^{\frac{1}{n+1}} a_1^2 \left(z - \frac{b_m}{a_1}\right) \left(z - \frac{b_{n+1}}{a_1}\right) & (n = 2m-1) \\ t^{\frac{1}{n+1}} a_1^2 \left(z - \frac{a_m}{a_1}\right) \left(z - \frac{b_{n+1}}{a_1}\right) & (n = 2m) \end{cases}.$$

T_4 : q -LUC system

Lemma

The transformation T_4 act on (φ_i) as

$$T_4(\varphi_{2i}) = \frac{1}{\alpha_{2i+2} \alpha_{2i+3}} \frac{(1 + \varphi_{2i+1})(\alpha_{2i+3} + \varphi_{2i+3})}{(1 + \varphi_{2i+3})(\alpha_{2i+1} + \varphi_{2i+1})} \varphi_{2i+2},$$

$$T_4(\varphi_{2i+1}) = \alpha_{2i+1} \alpha_{2i+2} \frac{\{1 + T_4(\varphi_{2i+2})\} \left\{ \frac{1}{\alpha_{2i+1} \alpha_{2i+2} \alpha_{2i+3}} + T_4(\varphi_{2i}) \right\}}{\{1 + T_4(\varphi_{2i})\} \left\{ \frac{1}{\alpha_{2i+3} \alpha_{2i+4} \alpha_{2i+5}} + T_4(\varphi_{2i+2}) \right\}} \varphi_{2i+3}.$$

Remark

The above system is not suitable as a q -Painlevé system because T_4 act on the parameters as

$$T_4(\alpha_{2i}) = \frac{1}{\alpha_{2i+1} \alpha_{2i+2} \alpha_{2i+3}}, \quad T_4(\alpha_{2i+1}) = \alpha_{2i+1} \alpha_{2i+2} \alpha_{2i+3} \alpha_{2i+4} \alpha_{2i+5}.$$

T_4 : q -LUC system

Let $t = (t_0, t_1)$ be a 2-tuples of independent variables defined by

$$t_0 = \prod_{i=0}^n \frac{1}{\varphi_{2i}} = \frac{1}{\beta^{\frac{1}{2}} (\beta')^{\frac{1}{2}}} \prod_{i=0}^n \frac{1}{\alpha_{2i}^{\frac{1}{2}}}, \quad t_1 = \prod_{i=0}^n \frac{\alpha_{2i+1}}{\varphi_{2i+1}} = \frac{\beta^{\frac{1}{2}}}{(\beta')^{\frac{1}{2}}} \prod_{i=0}^n \alpha_{2i}^{\frac{1}{2}} \alpha_{2i+1}.$$

Also let $c = (c_0, \dots, c_{2n+1})$ be a tuple of parameters defined by

$$c_{2i} = t_0^{\frac{1}{n+1}} t_1^{\frac{1}{n+1}} \alpha_{2i}, \quad c_{2i+1} = t_0^{-\frac{1}{n+1}} t_1^{-\frac{1}{n+1}} \alpha_{2i+1}.$$

Then we have

$$T_4(t_0) = q t_0, \quad T_4(t_1) = q t_1,$$

and

$$T_4^m(c_{2i}) = c_{2i}, \quad T_4^m(c_{2i+1}) = c_{2i+1} \quad (n = 2m - 1),$$

$$T_4^{2m+1}(c_{2i}) = c_{2i}, \quad T_4^{2m+1}(c_{2i+1}) = c_{2i+1} \quad (n = 2m).$$

We now regard each φ_i as a dependent variable on (t, c) and denote it by $\varphi_i(t; c)$.

T_4 : q -LUC system

With the aid of the action of T_4 on c , we define dependent variables and parameters on a lattice by

$$f_{i,n-i+2k}(t) = t_0^{\frac{1}{n+1}} \varphi_{2i}(t; T_4^{-k}(c)), \quad g_{i,n-i+2k-1}(t) = \frac{t_0^{-\frac{1}{n+1}} \varphi_{2i+1}(t; T_4^{-k}(c))}{T_4^{-k}(c_{2i+1})},$$

$$c_{i,n-i+2k} = \frac{T_4^{-k}(c_{2i})}{q^{\frac{1}{n+1}}}, \quad c_{i,n-i+2k-1} = \frac{1}{T_4^{-k}(c_{2i+1})},$$

where the indices i, j of $f_{i,j}(t), g_{i,j}(t), c_{i,j}$ are congruent modulo $n + 1$. Note that

$$\prod_{i=0}^n f_{i,n-i+2k} = 1, \quad \prod_{i=0}^n g_{i,n-i+2k-1} = 1 \quad (k = 0, \dots, n),$$

$$\frac{c_{i,n-i+2k} c_{i+1,n-i+2k+1}}{c_{i+1,n-i+2k} c_{i,n-i+2k+1}} = 1, \quad \frac{c_{i,n-i+2k-1} c_{i+1,n-i+2k}}{c_{i+1,n-i+2k-1} c_{i,n-i+2k}} = 1 \quad (i, k = 0, \dots, n).$$

We also set

$$\alpha = -t_1^{\frac{1}{n+1}}, \quad \beta = -t_0^{-\frac{1}{n+1}}, \quad \gamma = -t_0^{\frac{1}{n+1}}, \quad \delta = -t_1^{-\frac{1}{n+1}}.$$

T_4 : q -LUC system

Theorem ([Okubo-S 18])

The dependent variables $f_{i,n-i+2k}, g_{i,n-i+2k-1}$ satisfy

$$\begin{aligned} \frac{\bar{f}_{i,n-i+2k}}{f_{i+1,n-i+2k+1}} &= \frac{c_{i,n-i+2k}}{c_{i,n-i+2k+1}} \frac{(g_{i+1,n-i+2k} - \alpha)(g_{i,n-i+2k+1} - c_{i,n-i+2k+1}\beta)}{(g_{i,n-i+2k+1} - \alpha)(g_{i+1,n-i+2k} - c_{i+1,n-i+2k}\beta)}, \\ \frac{\bar{g}_{i,n-i+2k-1}}{g_{i+1,n-i+2k}} &= \frac{c_{i+1,n-i+2k-1}}{c_{i+1,n-i+2k}} \frac{(\bar{f}_{i+1,n-i+2k-1} - \bar{\gamma})(\bar{f}_{i,n-i+2k} - c_{i,n-i+2k}\delta)}{(\bar{f}_{i,n-i+2k} - \bar{\gamma})(\bar{f}_{i+1,n-i+2k-1} - c_{i+1,n-i+2k-1}\delta)}, \end{aligned}$$

where

$$\begin{aligned} \bar{f}_{i,n-i+2k} &= f_{i,n-i+2k}(q t), \quad \bar{g}_{i,n-i+2k-1} = g_{i,n-i+2k-1}(q t), \\ \bar{\alpha} &= q^{\frac{1}{n+1}} \alpha, \quad \bar{\beta} = q^{-\frac{1}{n+1}} \beta, \quad \bar{\gamma} = q^{\frac{1}{n+1}} \gamma, \quad \bar{\delta} = q^{-\frac{1}{n+1}} \delta. \end{aligned}$$

1 Introduction

2 Extended affine Weyl group

3 Higher order q -Painlevé system

4 Conclusion

In this talk we obtain the following object from cluster mutations.

- A birational representation of $\widetilde{W}((A_{2n+1} + A_1 + A_1)^{(1)})$
- The known three q -Painlevé systems
 - q -DS system $q\text{-}P_{(n+1,n+1)}$
 - Sakai's q -Garnier system
 - Tsuda's q -LUC system

We want to investigate the following properties of the above q -Painlevé systems from a viewpoint of the Kac-Moody Weyl group or the cluster algebra.

- τ -Functions
- Lax pairs
- Hypergeometric solutions

The following objects are also our next targets.

- Lax pair of q -Sasano system
- q -Painlevé systems of type $(A_{l-1} + A_{m-1} + A_{n-1})^{(1)}$

Thank you for your attention.